

- Non-zero particle masses are $m_t = 173$ GeV, $m_H = 125$ GeV, $m_Z = 91.188$ GeV, $m_W = 80.419$ GeV. The bottom-quark mass is set to $m_b = 4.7$ GeV in the four-flavour-scheme (4FS) simulations, and to $m_b = 0$ in the five-flavour-scheme (5FS) ones. The CKM matrix is $V_{\text{CKM}} = \mathbb{1}$, and the fine-structure constant is $\alpha = 1/132.507$.
- Renormalisation and factorisation scales are chosen as $\mu_R = \mu_F = \frac{1}{2} \sum_k m_T^{(k)}$, $m_T^{(k)}$ being the transverse mass of the k -th final-state particle. Independent variation of μ_R and μ_F in the range $[1/2, 2]$ is obtained in an exact way without rerunning the code, through the reweighting technique described in [1]. The uncertainty associated with this variation is shown as a dark band in the plots of the following sections.
- As parton distribution functions (PDFs) we use the MSTW 2008 NLO (68% c.l.) sets [2], relevant to four or five active flavours, depending on the flavour scheme employed in the simulation. PDF uncertainties are estimated according to the asymmetric-hessian prescription provided by the PDF set, and obtained automatically as explained in [1]. They are shown as a light band in the plots of the following sections. The value and the running of the strong coupling constant α_s are as well set according to the PDF set.
- Whenever relevant, jets are clustered by means of the anti- k_T algorithm [3] as implemented in FASTJET [4], with $\Delta R = 0.7$, and $p_T(j) > 20$ GeV (if not otherwise specified).
- Whenever relevant, photons are isolated by means of the Frixione smooth-cone criterion [5], with parameters $R_0 = 0.4$, $p_T(\gamma) > 20$ GeV, $\epsilon_\gamma = n = 1$.
- Loop-induced processes are generated in the 4FS.

Process	$\sigma_{\text{NLO}}(8 \text{ TeV})$ [fb]	$\sigma_{\text{NLO}}(100 \text{ TeV})$ [fb]	ρ
$pp \rightarrow H(m_t, m_b)$	$1.44 \cdot 10^4$ $^{+20\% +1\%}_{-16\% -2\%}$	$5.46 \cdot 10^{+5}$ $^{+28\% +2\%}_{-27\% -2\%}$	38
$gg \rightarrow H$ (LI)	$6.85 \cdot 10^{+3}$ $^{+33\% +1\%}_{-24\% -1\%}$	$2.21 \cdot 10^{+5}$ $^{+58\% +1\%}_{-39\% -1\%}$	32
$pp \rightarrow Hjj$ (VBF)	$1.61 \cdot 10^3$ $^{+1\% +2\%}_{-0\% -2\%}$	$7.40 \cdot 10^4$ $^{+3\% +2\%}_{-2\% -1\%}$	46
$pp \rightarrow Hjj$ (LI)	$1.89 \cdot 10^{+3}$ $^{+67\% +1\%}_{-37\% -1\%}$	$2.02 \cdot 10^{+5}$ $^{+66\% +0\%}_{-38\% -1\%}$	107
$pp \rightarrow Ht\bar{t}$	$1.21 \cdot 10^2$ $^{+5\% +3\%}_{-9\% -3\%}$	$3.25 \cdot 10^4$ $^{+7\% +1\%}_{-8\% -1\%}$	269
$pp \rightarrow Hb\bar{b}$ (4FS)	$2.37 \cdot 10^2$ $^{+9\% +2\%}_{-9\% -2\%}$	$1.21 \cdot 10^4$ $^{+2\% +2\%}_{-10\% -2\%}$	51
$pp \rightarrow Htj$	$2.07 \cdot 10^1$ $^{+2\% +2\%}_{-1\% -2\%}$	$5.21 \cdot 10^3$ $^{+3\% +1\%}_{-5\% -1\%}$	252
$pp \rightarrow HW^\pm$	$7.31 \cdot 10^2$ $^{+2\% +2\%}_{-1\% -2\%}$	$1.54 \cdot 10^4$ $^{+5\% +2\%}_{-8\% -2\%}$	21
$pp \rightarrow HZ$	$3.87 \cdot 10^2$ $^{+2\% +2\%}_{-1\% -2\%}$	$8.82 \cdot 10^3$ $^{+4\% +2\%}_{-8\% -2\%}$	23
$gg \rightarrow HZ$ (LI)	22.9 $^{+33\% +1\%}_{-23\% -2\%}$	$2.50 \cdot 10^{+3}$ $^{+35\% +1\%}_{-26\% -1\%}$	109
$pp \rightarrow HW^+W^-$ (4FS)	$4.62 \cdot 10^0$ $^{+3\% +2\%}_{-2\% -2\%}$	$1.68 \cdot 10^2$ $^{+5\% +2\%}_{-6\% -1\%}$	36
$gg \rightarrow HW^+W^-$ (LI)	N/A	N/A	[-]
$pp \rightarrow HZW^\pm$	$2.17 \cdot 10^0$ $^{+4\% +2\%}_{-4\% -2\%}$	$9.94 \cdot 10^1$ $^{+6\% +2\%}_{-7\% -1\%}$	46
$pp \rightarrow HW^\pm\gamma$	$2.36 \cdot 10^0$ $^{+3\% +2\%}_{-3\% -2\%}$	$7.75 \cdot 10^1$ $^{+7\% +2\%}_{-8\% -1\%}$	33
$pp \rightarrow HZ\gamma$	$1.54 \cdot 10^0$ $^{+3\% +2\%}_{-2\% -2\%}$	$4.29 \cdot 10^1$ $^{+5\% +2\%}_{-7\% -2\%}$	28
$gg \rightarrow HZ\gamma$ (LI)	$2.12 \cdot 10^{-3}$ $^{+34\% +1\%}_{-23\% -2\%}$	0.279 $^{+33\% +0\%}_{-25\% -1\%}$	132
$gg \rightarrow HH$ (LI)	5.46 $^{+34\% +2\%}_{-24\% -2\%}$	$7.74 \cdot 10^{+2}$ $^{+32\% +0\%}_{-24\% -1\%}$	142
$pp \rightarrow ZZ$	$7.18 \cdot 10^{+3}$ $^{+3\% +2\%}_{-3\% -2\%}$	$1.62 \cdot 10^{+5}$ $^{+6\% +2\%}_{-10\% -2\%}$	23
$gg \rightarrow ZZ$ (LI)	$4.93 \cdot 10^{+2}$ $^{+30\% +1\%}_{-21\% -1\%}$	$2.92 \cdot 10^{+4}$ $^{+42\% +1\%}_{-30\% -1\%}$	59
$pp \rightarrow Z\gamma$	$3.63 \cdot 10^{+4}$ $^{+5\% +2\%}_{-5\% -1\%}$	$6.00 \cdot 10^{+5}$ $^{+13\% +3\%}_{-16\% -3\%}$	17
$gg \rightarrow Z\gamma$ (LI)	$3.98 \cdot 10^{+2}$ $^{+29\% +1\%}_{-21\% -1\%}$	$1.70 \cdot 10^{+4}$ $^{+52\% +1\%}_{-35\% -1\%}$	43
$pp \rightarrow \gamma\gamma$	$1.15 \cdot 10^{+5}$ $^{+14\% +2\%}_{-15\% -1\%}$	$1.52 \cdot 10^{+6}$ $^{+32\% +8\%}_{-29\% -6\%}$	13
$gg \rightarrow \gamma\gamma$ (LI)	$2.54 \cdot 10^{+4}$ $^{+56\% +1\%}_{-37\% -1\%}$	$4.59 \cdot 10^{+5}$ $^{+89\% +3\%}_{-50\% -3\%}$	18
$pp \rightarrow HZZ$	$1.10 \cdot 10^0$ $^{+2\% +2\%}_{-2\% -2\%}$	$4.20 \cdot 10^1$ $^{+4\% +2\%}_{-6\% -1\%}$	38
$pp \rightarrow W^+W^-$ (4FS)	$5.30 \cdot 10^4$ $^{+3\% +2\%}_{-3\% -2\%}$	$1.12 \cdot 10^6$ $^{+9\% +2\%}_{-11\% -2\%}$	21
$gg \rightarrow W^+W^-$ (LI)	N/A	N/A	[-]
$pp \rightarrow HZZ$	$1.10 \cdot 10^0$ $^{+2\% +2\%}_{-2\% -2\%}$	$4.20 \cdot 10^1$ $^{+4\% +2\%}_{-6\% -1\%}$	38
$gg \rightarrow HZZ$ (LI)	$3.56 \cdot 10^{-2}$ $^{+35\% +2\%}_{-24\% -2\%}$	7.29 $^{+28\% +0\%}_{-22\% -1\%}$	205
$gg \rightarrow H\gamma\gamma$ (LI)	$2.24 \cdot 10^{-3}$ $^{+35\% +2\%}_{-24\% -2\%}$	0.374 $^{+31\% +0\%}_{-24\% -1\%}$	167
$pp \rightarrow HW^\pm j$	$3.18 \cdot 10^2$ $^{+4\% +2\%}_{-4\% -1\%}$	$1.07 \cdot 10^4$ $^{+2\% +2\%}_{-7\% -1\%}$	34
$pp \rightarrow HW^\pm jj$	$6.06 \cdot 10^1$ $^{+6\% +1\%}_{-8\% -1\%}$	$4.90 \cdot 10^3$ $^{+2\% +1\%}_{-6\% -1\%}$	81
$pp \rightarrow HZj$	$1.71 \cdot 10^2$ $^{+4\% +1\%}_{-4\% -1\%}$	$6.31 \cdot 10^3$ $^{+2\% +2\%}_{-7\% -1\%}$	37
$pp \rightarrow HZjj$	$3.50 \cdot 10^1$ $^{+7\% +1\%}_{-10\% -1\%}$	$2.81 \cdot 10^3$ $^{+2\% +1\%}_{-5\% -1\%}$	80

Table 1: Production of a single Higgs boson at the LHC and at a 100 TeV FCC-hh. The rightmost column reports the ratio ρ of the FCC-hh to the LHC cross sections. Theoretical uncertainties are due to scale and PDF variations, respectively. Monte-Carlo-integration error is always smaller than theoretical uncertainties, and is not shown. For $pp \rightarrow HVjj$, on top of the transverse-momentum cut of section ??, I require $m(j_1, j_2) > 100$ GeV, j_1 and j_2 being the hardest and next-to-hardest jets, respectively. Processes $pp \rightarrow Htj$ and $pp \rightarrow Hjj$ (VBF) do not feature jet cuts.

	Process	$\sigma_{\text{NLO}}(8 \text{ TeV}) [\text{fb}]$	$\sigma_{\text{NLO}}(100 \text{ TeV}) [\text{fb}]$	ρ
$pp \rightarrow$	$W^+W^-W^\pm$ (4FS)	$8.73 \cdot 10^1 \begin{smallmatrix} +6\% & +2\% \\ -4\% & -2\% \end{smallmatrix}$	$4.25 \cdot 10^3 \begin{smallmatrix} +9\% & +1\% \\ -9\% & -1\% \end{smallmatrix}$	49
$pp \rightarrow$	W^+W^-Z (4FS)	$6.41 \cdot 10^1 \begin{smallmatrix} +7\% & +2\% \\ -5\% & -2\% \end{smallmatrix}$	$4.01 \cdot 10^3 \begin{smallmatrix} +9\% & +1\% \\ -9\% & -1\% \end{smallmatrix}$	63
$gg \rightarrow$	W^+W^-Z (LI)	N/A	N/A	[-]
$pp \rightarrow$	$Z\gamma\gamma$	$6.06 \cdot 10^{+1} \begin{smallmatrix} +4\% & +2\% \\ -4\% & -1\% \end{smallmatrix}$	$1.37 \cdot 10^{+3} \begin{smallmatrix} +12\% & +2\% \\ -13\% & -2\% \end{smallmatrix}$	23
$gg \rightarrow$	$Z\gamma\gamma$ (LI)	$5.58 \cdot 10^{-2} \begin{smallmatrix} +28\% & +1\% \\ -21\% & -1\% \end{smallmatrix}$	$3.42 \begin{smallmatrix} +44\% & +1\% \\ -31\% & -1\% \end{smallmatrix}$	61
$pp \rightarrow$	$\gamma W^\pm Z$	$7.11 \cdot 10^1 \begin{smallmatrix} +8\% & +2\% \\ -7\% & -1\% \end{smallmatrix}$	$3.61 \cdot 10^3 \begin{smallmatrix} +12\% & +1\% \\ -12\% & -1\% \end{smallmatrix}$	51
$pp \rightarrow$	$W^\pm ZZ$	$2.16 \cdot 10^1 \begin{smallmatrix} +7\% & +2\% \\ -6\% & -2\% \end{smallmatrix}$	$1.36 \cdot 10^3 \begin{smallmatrix} +10\% & +1\% \\ -10\% & -1\% \end{smallmatrix}$	63
$pp \rightarrow$	γZZ	$2.24 \cdot 10^1 \begin{smallmatrix} +4\% & +2\% \\ -3\% & -2\% \end{smallmatrix}$	$6.62 \cdot 10^2 \begin{smallmatrix} +8\% & +2\% \\ -9\% & -1\% \end{smallmatrix}$	30
$gg \rightarrow$	γZZ (LI)	$1.13 \cdot 10^{-3} \begin{smallmatrix} +33\% & +1\% \\ -23\% & -2\% \end{smallmatrix}$	$0.13 \begin{smallmatrix} +34\% & +1\% \\ -25\% & -1\% \end{smallmatrix}$	115
$pp \rightarrow$	ZZZ	$5.97 \cdot 10^0 \begin{smallmatrix} +3\% & +2\% \\ -3\% & -2\% \end{smallmatrix}$	$2.55 \cdot 10^2 \begin{smallmatrix} +5\% & +2\% \\ -7\% & -1\% \end{smallmatrix}$	43
$gg \rightarrow$	ZZZ (LI)	N/A	N/A	[-]
$gg \rightarrow$	HHH (LI)	N/A	N/A	[-]
$gg \rightarrow$	$HHHH$ (LI)	$2.63 \cdot 10^{-5} \begin{smallmatrix} +39\% & +3\% \\ -26\% & -3\% \end{smallmatrix}$	$1.30 \cdot 10^{-2} \begin{smallmatrix} +23\% & +1\% \\ -18\% & -1\% \end{smallmatrix}$	494
$pp \rightarrow$	HHZ	$1.31 \cdot 10^{-1} \begin{smallmatrix} +2\% & +2\% \\ -2\% & -2\% \end{smallmatrix}$	$5.36 \cdot 10^0 \begin{smallmatrix} +3\% & +2\% \\ -5\% & -1\% \end{smallmatrix}$	41
$gg \rightarrow$	HHZ (LI)	$1.60 \cdot 10^{-2} \begin{smallmatrix} +36\% & +2\% \\ -24\% & -2\% \end{smallmatrix}$	$3.35 \begin{smallmatrix} +29\% & +0\% \\ -22\% & -1\% \end{smallmatrix}$	209
$gg \rightarrow$	$HH\gamma$ (LI)	N/A	N/A	[-]
$pp \rightarrow$	$W^+W^-W^\pm\gamma$ (4FS)	$6.78 \cdot 10^{-1} \begin{smallmatrix} +8\% & +2\% \\ -6\% & -2\% \end{smallmatrix}$	$7.42 \cdot 10^1 \begin{smallmatrix} +8\% & +1\% \\ -8\% & -1\% \end{smallmatrix}$	109
$pp \rightarrow$	$W^+W^-W^\pm Z$ (4FS)	$3.48 \cdot 10^{-1} \begin{smallmatrix} +8\% & +2\% \\ -7\% & -2\% \end{smallmatrix}$	$5.95 \cdot 10^1 \begin{smallmatrix} +7\% & +1\% \\ -7\% & -1\% \end{smallmatrix}$	171
$pp \rightarrow$	$W^+W^-W^+W^-$ (4FS)	$3.01 \cdot 10^{-1} \begin{smallmatrix} +7\% & +2\% \\ -6\% & -2\% \end{smallmatrix}$	$4.11 \cdot 10^1 \begin{smallmatrix} +7\% & +1\% \\ -6\% & -1\% \end{smallmatrix}$	137
$pp \rightarrow$	W^+W^-ZZ (4FS)	$2.01 \cdot 10^{-1} \begin{smallmatrix} +7\% & +2\% \\ -6\% & -2\% \end{smallmatrix}$	$3.34 \cdot 10^1 \begin{smallmatrix} +6\% & +1\% \\ -6\% & -1\% \end{smallmatrix}$	166
$pp \rightarrow$	$W^\pm ZZZ$	$3.40 \cdot 10^{-2} \begin{smallmatrix} +10\% & +2\% \\ -8\% & -2\% \end{smallmatrix}$	$7.06 \cdot 10^0 \begin{smallmatrix} +8\% & +1\% \\ -7\% & -1\% \end{smallmatrix}$	208
$pp \rightarrow$	$ZZZZ$	$8.72 \cdot 10^{-3} \begin{smallmatrix} +4\% & +3\% \\ -4\% & -2\% \end{smallmatrix}$	$8.05 \cdot 10^{-1} \begin{smallmatrix} +4\% & +2\% \\ -4\% & -1\% \end{smallmatrix}$	92
$pp \rightarrow$	$W^+W^-W^+W^-\gamma$ (4FS)	$5.18 \cdot 10^{-3} \begin{smallmatrix} +8\% & +3\% \\ -7\% & -2\% \end{smallmatrix}$	$1.58 \cdot 10^0 \begin{smallmatrix} +6\% & +1\% \\ -5\% & -1\% \end{smallmatrix}$	305
$pp \rightarrow$	$ZZZZZ$	$1.07 \cdot 10^{-5} \begin{smallmatrix} +5\% & +3\% \\ -4\% & -2\% \end{smallmatrix}$	$2.04 \cdot 10^{-3} \begin{smallmatrix} +3\% & +2\% \\ -3\% & -1\% \end{smallmatrix}$	191

Table 2: Production of multiple vector bosons at the LHC and at a 100 TeV FCC-hh. The rightmost column reports the ratio ρ of the FCC-hh to the LHC cross sections. Theoretical uncertainties are due to scale and PDF variations, respectively. Monte-Carlo-integration error is always smaller than theoretical uncertainties, and is not shown.

	Process	$\sigma_{\text{NLO}}(8 \text{ TeV})$ [fb]	$\sigma_{\text{NLO}}(100 \text{ TeV})$ [fb]	ρ
$pp \rightarrow$	$t\bar{t}\gamma$	$6.50 \cdot 10^2$ $^{+12\%}_{-13\%}$ $^{+2\%}_{-2\%}$	$1.24 \cdot 10^5$ $^{+11\%}_{-11\%}$ $^{+1\%}_{-1\%}$	192
$pp \rightarrow$	$t\bar{t}Z$	$1.99 \cdot 10^2$ $^{+10\%}_{-12\%}$ $^{+3\%}_{-3\%}$	$5.63 \cdot 10^4$ $^{+9\%}_{-10\%}$ $^{+1\%}_{-1\%}$	282
$pp \rightarrow$	$t\bar{t}W^\pm$	$2.05 \cdot 10^2$ $^{+9\%}_{-10\%}$ $^{+2\%}_{-2\%}$	$1.68 \cdot 10^4$ $^{+18\%}_{-16\%}$ $^{+1\%}_{-1\%}$	82
$pp \rightarrow$	$t\bar{t}\gamma j$	$1.22 \cdot 10^2$ $^{+17\%}_{-18\%}$ $^{+3\%}_{-3\%}$	$6.07 \cdot 10^4$ $^{+8\%}_{-10\%}$ $^{+1\%}_{-1\%}$	498
$pp \rightarrow$	$t\bar{t}Zj$	$3.51 \cdot 10^1$ $^{+15\%}_{-18\%}$ $^{+4\%}_{-4\%}$	$2.77 \cdot 10^4$ $^{+7\%}_{-9\%}$ $^{+1\%}_{-1\%}$	789
$pp \rightarrow$	$t\bar{t}W^\pm j$	$3.59 \cdot 10^1$ $^{+18\%}_{-18\%}$ $^{+2\%}_{-2\%}$	$1.36 \cdot 10^4$ $^{+14\%}_{-13\%}$ $^{+1\%}_{-1\%}$	379
$pp \rightarrow$	$t\bar{t}W^\pm jj$	$5.67 \cdot 10^0$ $^{+24\%}_{-23\%}$ $^{+3\%}_{-2\%}$	$6.52 \cdot 10^3$ $^{+11\%}_{-14\%}$ $^{+1\%}_{-1\%}$	1150
$pp \rightarrow$	$t\bar{t}W^+W^-$ (4FS)	$2.27 \cdot 10^0$ $^{+11\%}_{-13\%}$ $^{+3\%}_{-3\%}$	$1.10 \cdot 10^3$ $^{+9\%}_{-9\%}$ $^{+1\%}_{-1\%}$	486
$pp \rightarrow$	$t\bar{t}\gamma\gamma$	$2.23 \cdot 10^0$ $^{+14\%}_{-13\%}$ $^{+2\%}_{-1\%}$	$4.81 \cdot 10^2$ $^{+13\%}_{-11\%}$ $^{+1\%}_{-1\%}$	216
$pp \rightarrow$	$t\bar{t}Z\gamma$	$1.11 \cdot 10^0$ $^{+12\%}_{-13\%}$ $^{+2\%}_{-2\%}$	$4.20 \cdot 10^2$ $^{+10\%}_{-9\%}$ $^{+1\%}_{-1\%}$	378
$pp \rightarrow$	$t\bar{t}W^\pm Z$	$9.71 \cdot 10^{-1}$ $^{+10\%}_{-11\%}$ $^{+3\%}_{-2\%}$	$1.68 \cdot 10^2$ $^{+16\%}_{-13\%}$ $^{+1\%}_{-1\%}$	173
$pp \rightarrow$	$t\bar{t}ZZ$	$4.47 \cdot 10^{-1}$ $^{+8\%}_{-10\%}$ $^{+3\%}_{-2\%}$	$1.58 \cdot 10^2$ $^{+15\%}_{-12\%}$ $^{+1\%}_{-1\%}$	353

Table 3: Production of a top-antitop pair in association with up to two electroweak vector bosons, and with an electroweak boson and up to two jets, at the LHC and at a 100 TeV FCC-hh. The rightmost column reports the ratio ρ of the FCC-hh to the LHC cross sections. Processes $pp \rightarrow t\bar{t}Vj(j)$ feature a cut of $p_T(j) > 100$ GeV. Theoretical uncertainties are due to scale and PDF variations, respectively. Monte-Carlo-integration error is always smaller than theoretical uncertainties, and is not shown.

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